

ETM ENGINEERING AND PERFORMANCE OPTIMIZATION OF LEAD-FREE  $K_2GeBr_6$  PEROVSKITE SOLAR CELLSAbdul Mateen Arshad<sup>1</sup>, Azmat Ali<sup>2</sup>, Ahmed Salim<sup>3</sup>, Mohsin M. Tarar<sup>4</sup>, SU Rahman<sup>5</sup>, Raja Mohsin Ali Khan<sup>6</sup><sup>1,2,6</sup>Department of Physics, University of Chakwal, Chakwal, Pakistan<sup>3</sup>Department of Electrical Engineering, NAMAL University Mianwali, Mianwali, Pakistan<sup>4</sup>Department of Electronics Engineering, University of Chakwal, Chakwal, Pakistan<sup>5</sup>Department of Mathematics, University of Chakwal, Chakwal, Pakistan<sup>1</sup>mateenarshad765@gmail.com, <sup>2</sup>azmat.chakwal@gmail.com, <sup>3</sup>ahmed.salim@namal.edu.pk, <sup>4</sup>mohsin.tarar@uoc.edu.pk, <sup>5</sup>shafiq.urrhman@uoc.edu.pk, <sup>6</sup>mohsinkhan5252470@gmail.comDOI: <https://doi.org/10.5281/zenodo.19333001>**Keywords**

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**Abstract**

In this work, we examine a lead-free double perovskite absorber,  $K_2GeBr_6$ , in device architecture FTO/ETM/ $K_2GeBr_6$ / $WO_3$ /Au via SCAPS-1D simulations. A systematic comparison was conducted among three different electron transport materials (ETMs),  $TiO_2$ ,  $SnO_2$ , and  $WO_3$ , with  $WO_3$  as the hole transport material. The results indicate that  $TiO_2$  is the best interface with  $K_2GeBr_6$  due to device electrical parameters such as highest efficiency of 26.8%,  $V_{oc}=1.46$  V,  $J_{sc}=24.9$  mA/cm<sup>2</sup>, and FF=73.5%. Although  $SnO_2$  and  $WO_3$  shows promising characteristics, such as greater recombination losses with compromised performance. The results conform  $TiO_2$  as the most efficient ETM of  $K_2GeBr_6$  and confer its viability as the sustainable, nontoxic absorber material in the design of the future environmentally friendly perovskite.

**INTRODUCTION:**

The deepening global energy crisis, combined with the environmental degradation associated with the consumption of fossil fuels, has turned the technology of renewable energy into a prime focus of scientific research[1]. Among the available renewable energy sources, solar energy is the most easily accessible and has the potential to fulfill the energy requirements of the modern world in the long run[2, 3]. With time, the solar cell technology has evolved immensely from first-generation crystalline silicon solar cells to second-generation thin-film solar cells, and now to third-generation

materials. Although, silicon-based photovoltaics lead the market, their expensive processing and internal efficiency limits have initiated the search for alternative materials[4, 5]. In this regard, perovskite solar cells (PSCs) have transformed the field in a huge way, with a sudden spurt in efficiency from less than 4% in the early 2000s to more than 25% within two decades. This advance is due to the exceptional properties of perovskite absorbers, such as tunable bandgap, high optical absorption, long carrier diffusion lengths, and low-cost, simple fabrication processes [6-8].



Fig 1. Simulated diagram of  $K_2GeBr_6$  based perovskite solar cell

Though traditional PSCs have been developing at a fast pace yet they suffer from extreme issues concerning the usage of toxic lead. Despite being highly efficient but are environment- and health-risky owing to the possibility of lead leakage during device operation or waste[9, 10]. Further, these materials are typically plagued by poor thermal and chemical stability, limiting their large-scale use. To overcome these limitations, researchers

have focused on the field of lead-free perovskites, specifically double perovskite structures wherein the toxic  $Pb^{2+}$  is substituted with non-toxic species without sacrificing useful optoelectronic properties[11]. This approach not only removes toxicity but also enhances the device stability, thus, lead-free perovskites demonstrate a promising path towards green photovoltaics [12-14].

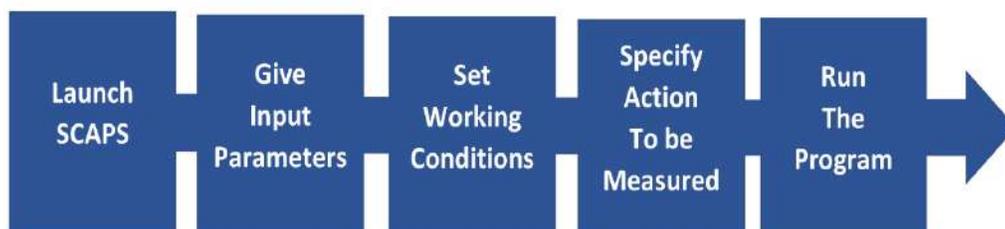


Fig 2. SCAPS working Procedure

Several classes of lead-free perovskites have been explored, including Bi-, Sb-, Sn-, and Cu-based materials. Antimony and bismuth-based double perovskites are stable but are plagued by indirect bandgaps and poor absorption that severely restrict the photocurrent generation. Perovskites

based on tin have appropriate bandgaps but are very unstable due to the oxidation of  $\text{Sn}^{2+}$  to  $\text{Sn}^{4+}$ , which adversely affects the performance. Copper-based alternatives, though green, have deep trap states and poor carrier lifetimes, which restrict practical efficiency[15, 16].

**Table 1: Material parameters of different ETMs, HTMs and Absorber.**

Material parameters	TiO <sub>2</sub> (ETM)	WO <sub>3</sub> (ETM)	SnO <sub>2</sub> (ETM)	MoO <sub>3</sub> (HTM)	K <sub>2</sub> GeBr <sub>6</sub> (Absorber)
Thickness(μm)	0.065	0.1	0.1	0.4	1.7
Band gap (eV)	3.2	2.6	2.6	3	1.6
Electron affinity (eV)	4.0	3.8	4.0	2.5	3.950
dielectric permittivity (relative)	10.0	4.8	9.0	12.5	3.420
CB effective density of states ( $\text{cm}^{-3}$ )	$1.0 \times 10^{21}$	$2.2 \times 10^{21}$	$2.2 \times 10^{18}$	$2.2 \times 10^{18}$	$1.0 \times 10^{17}$
VB effective density of states ( $\text{cm}^{-3}$ )	$2.0 \times 10^{20}$	$2.2 \times 10^{21}$	$1.8 \times 10^{19}$	$2.2 \times 10^{16}$	$1.0 \times 10^{18}$
electron thermal velocity ( $\text{cm s}^{-1}$ )	$1.0 \times 10^7$	$1.0 \times 10^7$	$1.0 \times 10^7$	$1.0 \times 10^7$	$1.0 \times 10^7$
hole thermal velocity ( $\text{cm s}^{-1}$ )	$1.0 \times 10^7$	$1.0 \times 10^7$	$1.0 \times 10^7$	$1.0 \times 10^7$	$1.0 \times 10^7$
electron mobility ( $\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$ )	$2.0 \times 10^1$	$3.0 \times 10^1$	$1.0 \times 10^2$	25	$1.6 \times 10^1$
hole mobility ( $\text{cm}^2 \text{V}^{-1} \text{s}^{-1}$ )	$1.0 \times 10^1$	$1.0 \times 10^1$	$2.5 \times 10^1$	100	$1.6 \times 10^1$
shallow uniform donor density ND (1/cm <sup>3</sup> )	$5.0 \times 10^{19}$	$6.35 \times 10^{17}$	$1.0 \times 10^{18}$	0	$1.0 \times 10^{17}$
shallow uniform acceptor density NA (1/cm <sup>3</sup> )	0	0	0	$1.0 \times 10^{18}$	$1.0 \times 10^{17}$

In the present work, K<sub>2</sub>GeBr<sub>6</sub> has been selected as the absorber owing to its advantageous optoelectronic and structural properties, as indicated by first-principles investigations concerning Ge-based double perovskites. Germanium based double Perovskites such that K<sub>2</sub>GeBr<sub>6</sub>/ K<sub>2</sub>GeI<sub>6</sub> have direct bandgap of 1.60/1.30 eV, which are suitable for absorption of light. These compounds also thermodynamically stable and exhibits stable structure validated by Goldschmidt tolerance factors[17]. Overall, the results of the theories verify K<sub>2</sub>GeBr<sub>6</sub> as optically favorable, structurally stable, and non-toxic

absorption material—and thus present an excellent candidate material for environmentally friendly high-efficiency solar cells[19, 20].

The electron hole pairs in the solar cell device is generated in absorber material by the incident of light which determines how efficient the device is? The absorber's bandgap, absorption spectrum, dielectric constants, optical conductivity, seebeck coefficient and defect tolerance govern how carriers are generated, how they recombine, and the efficiency of the solar cell. The electron transport material (ETM) plays a critical role in electron extraction, transport and in hole blocking

which reduces the recombination losses [21, 22]. The ETM is selected due to its compatibility with the absorber material. In this study three different ETMs ( $\text{TiO}_2$ ,  $\text{WO}_3$ , and  $\text{SnO}_2$ ) are used to determine their impact on device performance.  $\text{TiO}_2$  provides favorable conduction band alignment, suppresses recombination losses, and offers strong thermal and optical stability. Its compatibility with  $\text{K}_2\text{GeBr}_6$  absorbers ensures efficient carrier extraction and consistent device

performance. On the other  $\text{WO}_3$ , has a wide bandgap which support favorable band alignment and stable operation with temperatures. Its chemical stability shows long-term reliability.  $\text{SnO}_2$  is another promising ETM, having high electron mobility, optical transparency with low-temperature processing, which enable scalable fabrication while minimizing parasitic absorption and enhancing photocurrent generation.

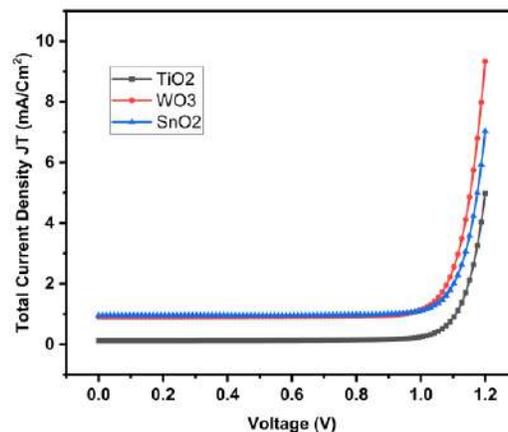


Figure 3: Total Current Density  $J$  ( $\text{mA}/\text{cm}^2$ ) vs Voltage  $V$  (V) for Optimized Solar cell with different ETMS

In this study  $\text{WO}_3$  is used as hole transport material (HTM) due to its valence band alignment with the  $\text{K}_2\text{GeBr}_6$  having efficient hole extraction and low recombination at the back interface. Its chemical stability and hardness shows the device as a whole for durability. Using  $\text{WO}_3$  as hole transport material with various electron transport material (ETM), it is feasible to make a numerical comparative analysis of the relative effect of  $\text{TiO}_2$ ,  $\text{WO}_3$ , and  $\text{SnO}_2$  based device performance. The Solar Cell Capacitance Simulator (SCAPS-1D) is used for numerical device simulations. The

technique enables detailed examination of generation, recombination, current-voltage characteristics, and quantum efficiency under various structural and defect conditions [5, 22]. The results show  $\text{K}_2\text{GeBr}_6$  with  $\text{TiO}_2$  offer excellent efficiency and stability compared to the devices with  $\text{WO}_3$  or  $\text{SnO}_2$  as ETMs[23, 24]. Study of  $\text{K}_2\text{GeBr}_6$ -based solar cells with various ETMs is an important step in the direction of highly-efficient, stable, eco-friendly, nontoxic perovskite photovoltaics.

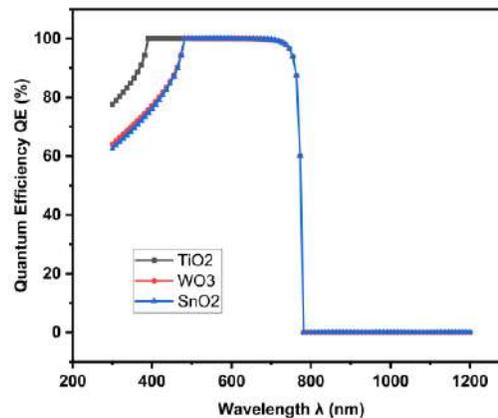


Figure 4: Quantum Efficiency QE (%) vs Wavelength ( $\lambda$ ) of optimized solar cell with different ETMs

The non-toxic absorber with well-matched transport material presents a potential replacement for conventional solar cell. This research adds to the emerging analysis and study on lead-free perovskites, positioning  $K_2GeBr_6$

with  $TiO_2$  ETM and  $WO_3$  HTM as excellent device articture. Material selection with device optimization provides scientific guidance and practical advice for realizing sustainable solar cell technologies.

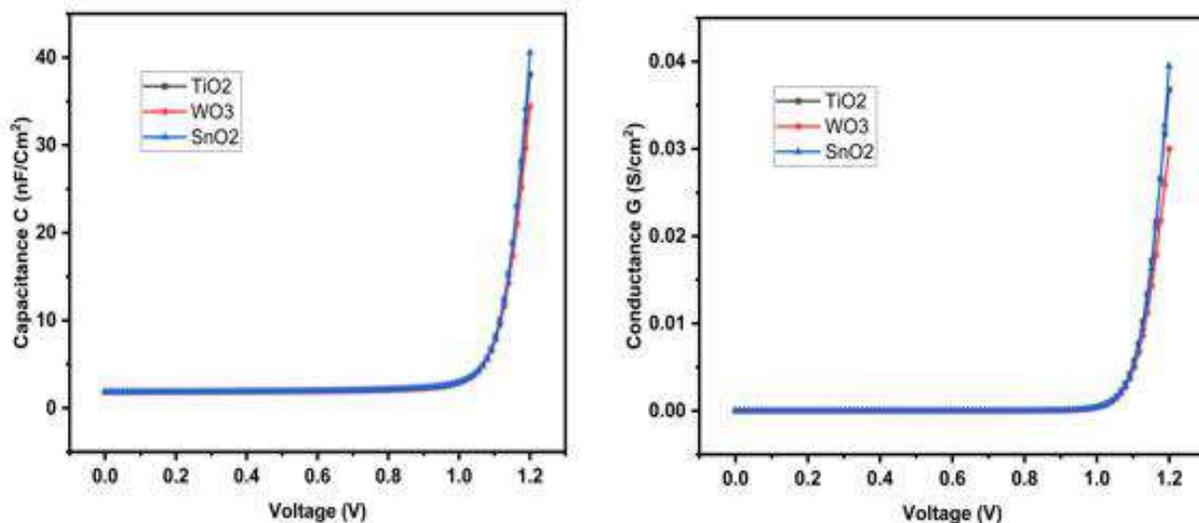


Figure 5: (a)-Capacitance and (b)-Conductance vs Voltage V(V) of optimized solar cell with different HTMs

## Methodology

### Numerical Simulations

The performance and effectiveness of solar cell were examined by using the SCAPS-1D (version 3.3.12) solar cell capacitance simulator, and critical device parameters were analyzed. This simulator is designed to model multi-layer solar cells and to calculated the electrical properties of device, such

as open-circuit voltage (VOC), short-circuit current density (JSC), band structure, quantum efficiency (QE%), and power conversion efficiency (PCE%). SCAPS-1D relies on solving partial differential equations (PDEs), including Poisson's equation and the continuity equations for both holes and electrons.[25, 26]

$$\dots \frac{\partial^2 \Psi}{\partial x^2} = -\frac{\partial E}{\partial x} = -\frac{\rho}{\epsilon_x} = -\frac{\rho}{\epsilon_x} [\mathbf{p} - \mathbf{n} + N_d^+ - N_a^- \pm N_{def}] \quad (1)$$

The symbol  $\Psi$  represents electrostatic potential, the term  $\epsilon_s$  represents relative permittivity, 'q' represents charge, 'n' represents electrons, and 'p' represents holes. ' $N_d^+$ ' denotes donor density, ' $N_a^-$ ' denotes acceptor density, and ' $N_{def}$ ' represents defect density for donor/acceptor. The equation describing continuity for holes and electrons is as follows:

$$\frac{\partial j_p}{\partial x} + G - U_p(n, p) = 0 \quad (2)$$

And

$$\frac{\partial j_n}{\partial x} + G - U_p(n, p) = 0 \quad (3)$$

The rates of hole recombination and electron recombination are represented by the variables  $U_p(n, p)$ ,  $U_n(n, p)$  respectively. The charge carrier generation rate, hole current density and electron current density is represented by the variable  $G$ ,  $J_p$  and  $J_n$ . The carrier current density can be obtained from these parameters:

$$j_p = qn\mu_p E - qD_p \frac{\partial p}{\partial x} \quad (4)$$

$$j_n = qn\mu_n E + qD_n \frac{\partial n}{\partial x} \quad (5)$$

The diffusion coefficients in this context are denoted as  $D_p$  and  $D_n$ , while the carrier mobilities are represented by  $\mu_n$  and  $\mu_p$ , with the charge denoted as  $q$ [5]. The program, developed at the University of Ghent[27], was introduced to the scientific community during the 2nd PV World Conference in Vienna for use in research and development.

#### Device Modeling and Analysis:

The simulated device is in the form of a planar single-junction structure made of FTO/ETM/ $K_2GeBr_6$ / $WO_3$ /Au, where  $K_2GeBr_6$  is the absorber,  $WO_3$  is the hole transport material (HTM), and the electron transport material (ETM) is  $TiO_2$ ,  $WO_3$ , or  $SnO_2$  to test their effect on the performance of the device. The front contact is made up of the transparent conducting oxide FTO, and Au is the back contact. Figure 1 depicts the general structure of the proposed device, showing the sequential layering of the elements and their respective functions in transporting and collecting charges.

Numerical computation was done using SCAPS-1D[28, 29], an open-access program to simulate the electronic and optical properties of thin-film solar cells. The simulation procedure and modeling process are shown in Figure 2, and the electronic and physical properties of all materials, including bandgap, electron affinity, permittivity,

mobility, doping, and defect density, are presented in Table 1. The research was conducted by comparing the performance of the three ETMs while holding the absorber and the HTM constant, hence determining their contribution to efficiency and stability improvement in  $K_2GeBr_6$ -based devices.[30]

#### Research Results and Discussions:

In Figure 3, At 0 V,  $SnO_2$  (0.95 mA/cm<sup>2</sup>) and  $WO_3$  (0.91 mA/cm<sup>2</sup>) both exhibit much higher current densities than  $TiO_2$  (0.12 mA/cm<sup>2</sup>), which demonstrates enhanced carrier extraction. At 1.2 V, the maximum current density is achieved by  $WO_3$  (9.33 mA/cm<sup>2</sup>), followed by  $SnO_2$  (7.02 mA/cm<sup>2</sup>) and  $TiO_2$  (4.98 mA/cm<sup>2</sup>). This suggests that  $SnO_2$  performs better at low bias due to its higher electron mobility, while  $WO_3$  performs better at high voltages due to favorable band alignment, with  $TiO_2$  consistently lagging behind due to recombination losses[31].

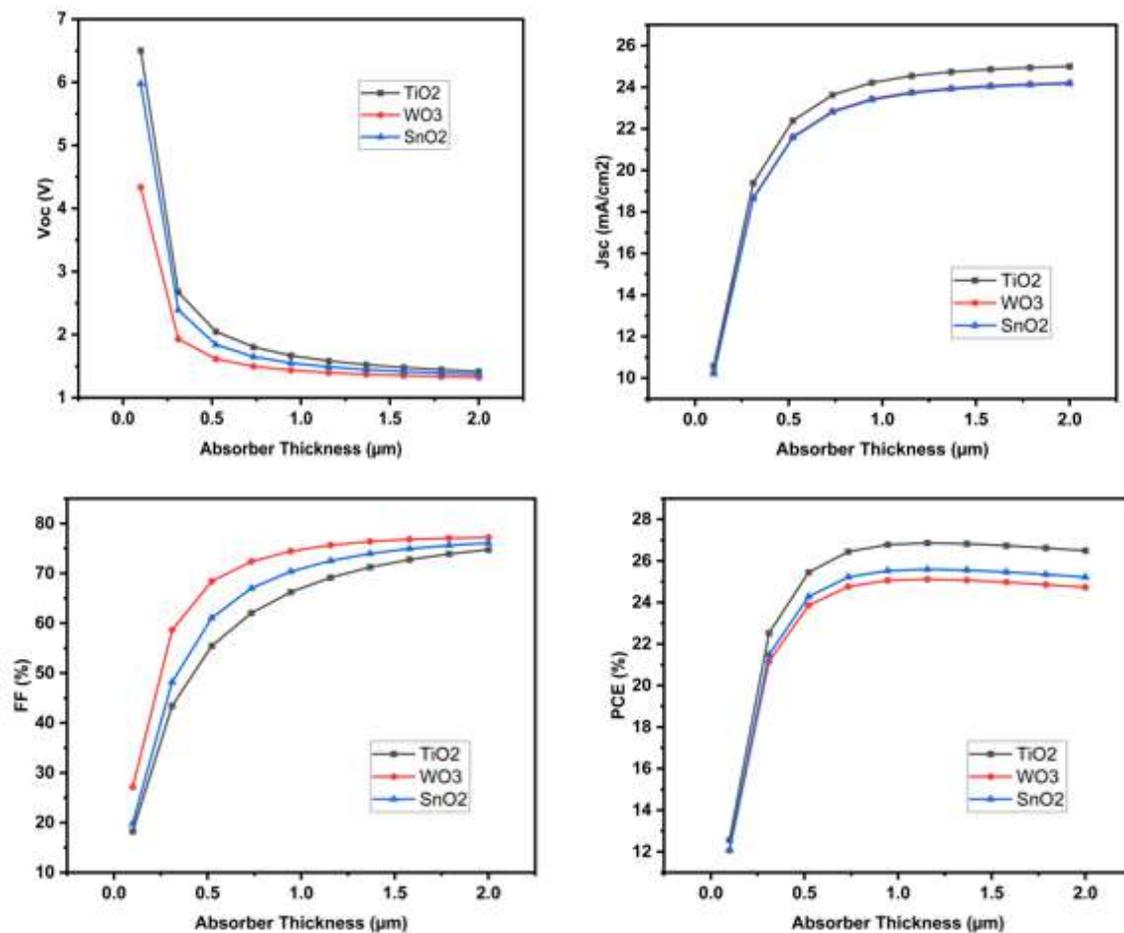
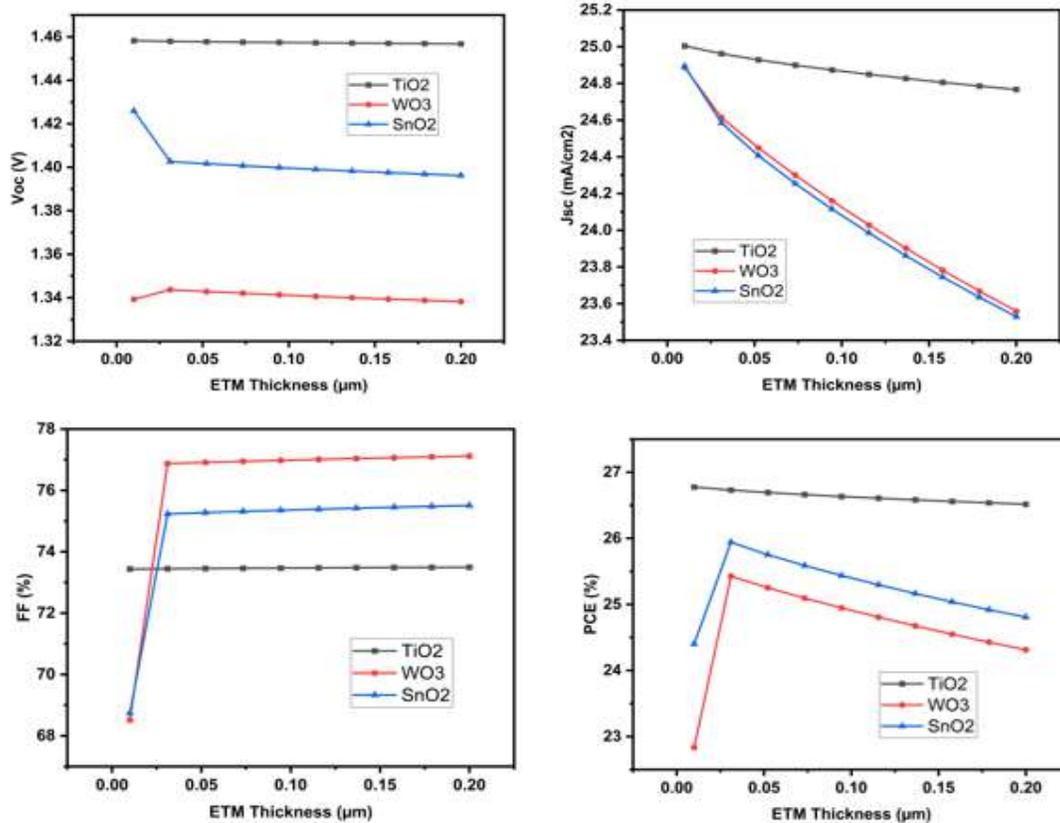


Figure 6: Variation of Electrical parameters vs Thickness of ( $K_2GeBr_6$ ) absorber layer with different ETMs

In Figure 4, QE increases from 300 nm with the highest  $TiO_2$  (77.7%) to  $WO_3$  (64.1%) and  $SnO_2$  (62.5%). At 390–400 nm,  $TiO_2$  is at 100%, whereas  $WO_3$  and  $SnO_2$  are 76–78%. In 480–600 nm, all devices are at about 100% QE, with perfect absorption, whereas above 700 nm QE drops for all.  $TiO_2$  performs better in the UV–visible range (300–500 nm) due to higher absorption edge, whereas  $WO_3$  and  $SnO_2$  are comparable in the mid-visible[32]. This verifies ETM selection primarily impacts QE at shorter wavelengths, such as comparable behavior in the flat absorption region

All the devices at 0 V possess nearly equal capacitance ( $TiO_2 = 1.83 \text{ nF/cm}^2$ ,  $WO_3 = 1.86 \text{ nF/cm}^2$ ,  $SnO_2 = 1.86 \text{ nF/cm}^2$ ), (Figure 5), and low

conductance ( $10^{-5} \text{ S/cm}^2$ )[33]. Capacitance and conductance increase slowly with increase in voltage up to 0.5 V, then very fast after 0.8 V, indicating improved storage and transport of charge on application of forward bias. The maximum capacitance ( $40.6 \text{ nF/cm}^2$ ) and conductance ( $3.95 \times 10^{-2} \text{ S/cm}^2$ ) at 1.2 V for  $SnO_2$  are followed by  $TiO_2$  ( $38.1 \text{ nF/cm}^2$ ,  $3.68 \times 10^{-2} \text{ S/cm}^2$ ), and the lowest capacitance ( $34.4 \text{ nF/cm}^2$ ) and conductance ( $3.0 \times 10^{-2} \text{ S/cm}^2$ ) are for  $WO_3$ . It means that  $SnO_2$  allows more efficient storage and transfer of charge, according to its higher mobility of electrons, whereas  $WO_3$  is slow due to less efficient carrier transport.  $TiO_2$  is average with consistent but average performance.



7: Variation of Electrical parameters vs Thickness of different ETMs

Figure 6 illustrates the influence of  $K_2GeBr_6$  absorber thickness (0.1–2 µm) on  $TiO_2$ ,  $WO_3$ , and  $SnO_2$  devices.  $V_{oc}$  decreases with greater thickness ( $TiO_2$ : 6.50 to 1.418 V,  $WO_3$ : 4.34 to 1.323 V,  $SnO_2$ : 5.98 to 1.372 V) due to greater recombination in thicker absorbers.  $J_{sc}$  increases very sharply in the beginning ( $TiO_2$ : 10.58 to 24.99 mA/cm<sup>2</sup>,  $WO_3$ : 10.25 to 24.21 mA/cm<sup>2</sup>,  $SnO_2$ : 10.21 to 24.17 mA/cm<sup>2</sup>) as more photons are absorbed and then becomes constant beyond 1.0 µm. FF increases ( $TiO_2$ : 18.26 to 74.74%,  $WO_3$ : 27.13 to 77.18%,  $SnO_2$ : 19.79 to 76.08%) as thicker absorbers enhance charge collection. PCE increases very sharply in the beginning ( $TiO_2$ : 12.56 to 26.49%,  $WO_3$ : 12.07 to 24.73%,  $SnO_2$ : 12.08 to 25.22%) and then remains constant, with  $TiO_2$  devices being the most efficient, followed by  $SnO_2$  and  $WO_3$ . This indicates that the optimal thickness of the absorber is a trade-off between photon absorption and recombination to achieve maximum efficiency[11].

Figure 7 illustrates the effect of ETM thickness (0.01–0.2 µm) on  $TiO_2$ ,  $WO_3$ , and  $SnO_2$   $K_2GeBr_6$  devices.  $V_{oc}$  decreases slightly with increased thickness ( $TiO_2$ : 1.458 to 1.457 V,  $WO_3$ : 1.339 to 1.338 V,  $SnO_2$ : 1.426 to 1.396 V) due to greater carrier transit length and recombination.  $J_{sc}$  decreases gradually too ( $TiO_2$ : 25.00 to 24.77 mA/cm<sup>2</sup>,  $WO_3$ : 24.89 to 23.56 mA/cm<sup>2</sup>,  $SnO_2$ : 24.89 to 23.53 mA/cm<sup>2</sup>) as thicker ETM retards electron extraction to a certain degree. FF increases slightly ( $TiO_2$ : 73.44 to 73.50%,  $WO_3$ : 68.51 to 77.12%,  $SnO_2$ : 68.73 to 75.51%) as series resistance decreases for thicker ETMs. As a result, PCE decreases slightly ( $TiO_2$ : 26.78 to 26.52%,  $WO_3$ : 22.84 to 24.31%,  $SnO_2$ : 24.40 to 24.81%), with the maximum efficiency for  $WO_3$  devices among intermediate thickness, then  $SnO_2$  and  $TiO_2$ [34]. This shows that ETM thickness needs to be optimized such that charge extraction and recombination are balanced for maximum efficiency.

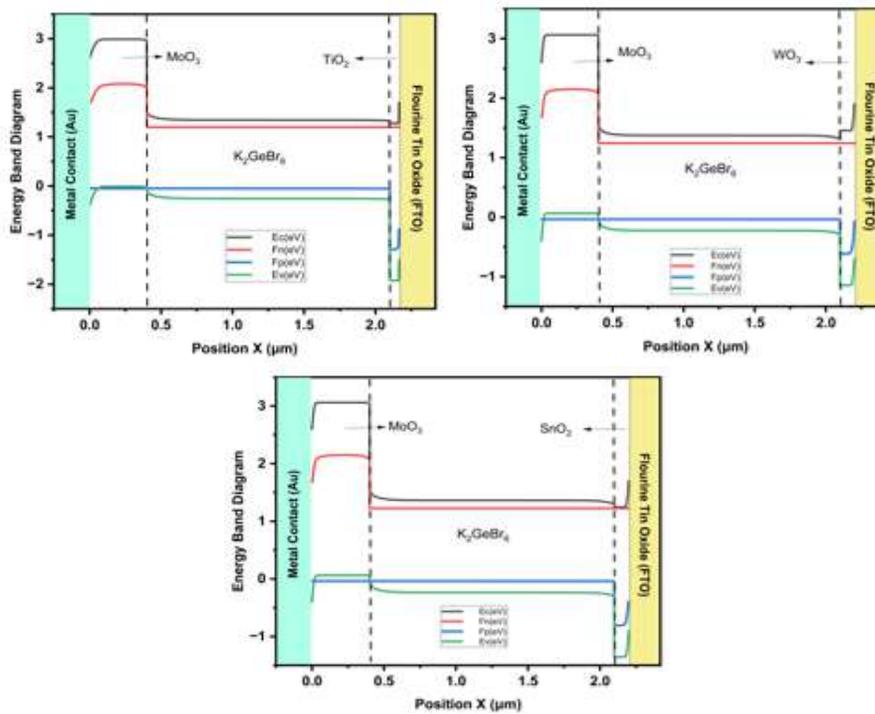


Figure 8: Energy band and layers thickness information diagram of optimized solar cell with different ETMs

In all three devices,  $E_c/E_v$  determine carrier transport and  $F_n/F_p$  splitting controls  $V_{oc}$ . (Figure 8)[35], For the  $TiO_2$  device,  $E_c$  at the ETL/absorber has a near-flat CBO ( $-2.0 \times 10^{-4}$  eV) and  $E_v$  at the absorber/HTM is also aligned, but the  $F_n-F_p$  separation in the absorber is 1.251 eV, the lowest, with stronger recombination and lower  $V_{oc}$ , as expected for  $TiO_2$ 's trap-related losses. For the  $WO_3$  device, offsets are still small (CBO  $\approx -6.2 \times 10^{-4}$  eV, VBO  $\approx -1.4 \times 10^{-4}$  eV), and  $F_n-F_p$  reaches 1.280 eV, the highest, with suppressed recombination and strongest  $V_{oc}$ . For  $SnO_2$ , offsets are similarly flat (CBO  $\approx -6.2 \times 10^{-4}$  eV, VBO  $\approx -1.2 \times 10^{-4}$  eV) but  $E_c$  bends sharply ( $-1.46 \times 10^{-2}$  eV/ $\mu m$ ), increasing the built-in field and extraction, resulting in  $F_n-F_p$  of 1.264 eV, somewhat lower than  $WO_3$  but higher than  $TiO_2$ , with stronger  $J_{sc}$  at little cost in  $V_{oc}$ . Efficiency potential thus follows  $WO_3 > SnO_2 > TiO_2$ .

Figure 9 shows that an increase in defects in the absorber severely degrades device performance as it leads to higher non-radiative recombination[36].  $V_{oc}$  is very high at low defect density ( $10^{13} \text{ cm}^{-3}$ ). (3.83 V for  $TiO_2$ , 3.02 V for  $SnO_2$ , 2.36 V for  $WO_3$ ) but drops to below 1.1 V at high density  $10^{17} \text{ cm}^{-3}$ .  $J_{sc}$  is unchanged at low density (24.9  $\text{ mA/cm}^2$ ) but drops to 21.4 for  $TiO_2$ , 20.4 for  $WO_3$ , and 20.5 for  $SnO_2$  at high density, indicating weaker collection of carriers. The FF increases at moderate concentrations of defects (75%) but then steeply drops to 39–45% at high density  $10^{17} \text{ cm}^{-3}$ . As a result, PCE drops from 29% for  $TiO_2$  and 28% for  $SnO_2$  and  $WO_3$  at high density  $10^{17} \text{ cm}^{-3}$ . Generally, higher defect density reduces  $V_{oc}$ ,  $J_{sc}$ , FF, and PCE due to higher trap-assisted recombination, although  $TiO_2$  shows a bit better  $V_{oc}$  and efficiency even at high concentrations.

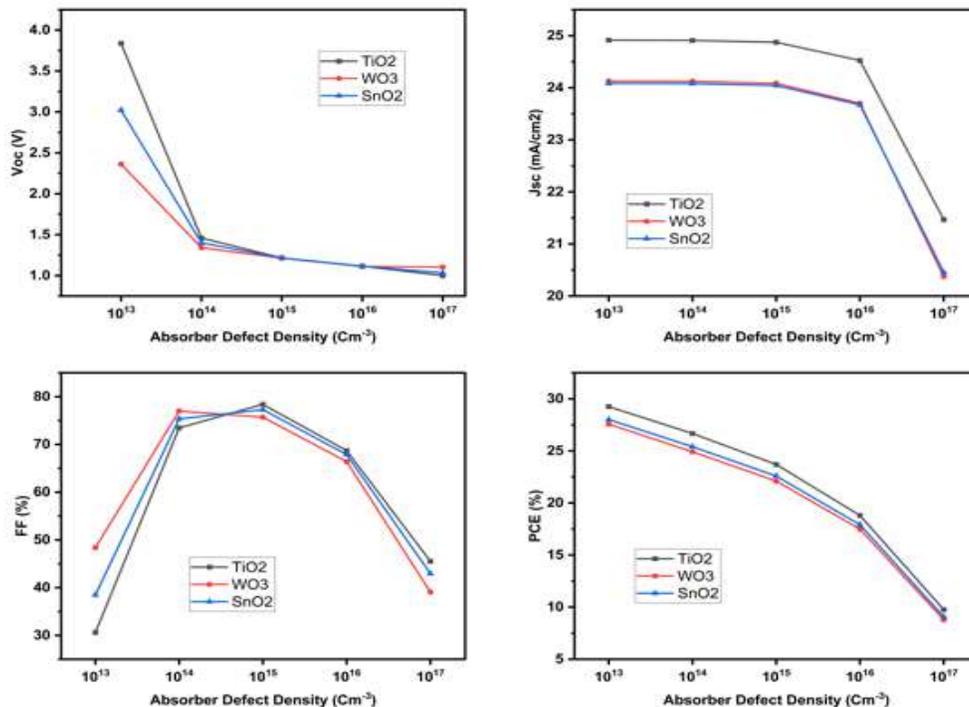


Figure 9: Variation of Electrical parameters vs Absorber Defect Density with different ETMs

In Figure 10, the depth profiles show that TiO<sub>2</sub> ETM yields more carrier generation ( $10^{21} \text{ cm}^{-3} \text{ s}^{-1}$  range) than WO<sub>3</sub> and SnO<sub>2</sub> ( $10^{21} \text{ cm}^{-3} \text{ s}^{-1}$  but consistently lower), reflecting stronger coupling of absorption and transport. While recombination for TiO<sub>2</sub> starts at about  $10^7 \text{ cm}^{-3} \text{ s}^{-1}$ , its peak ( $1.9 \times 10^{20} \text{ cm}^{-3} \text{ s}^{-1}$ ) is lower than WO<sub>3</sub>/SnO<sub>2</sub> ( $2.5 \times 10^{20} \text{ cm}^{-3} \text{ s}^{-1}$ ), indicating lower bulk losses. Likewise, SRH recombination in TiO<sub>2</sub> is lower in WO<sub>3</sub>/SnO<sub>2</sub> at depth, indicating effective trap-assisted decay suppression. Thus, while with marginally greater interfacial activity, TiO<sub>2</sub> gives the ideal balance of high generation and controlled recombination and is hence the most efficient ETM[37].

Figure 11 illustrates As CB-DOS is raised from  $1 \times 10^{16}$  to  $1 \times 10^{20} \text{ cm}^{-3}$ , Voc lowers (TiO<sub>2</sub>: 2.06 to 1.16 V, WO<sub>3</sub>: 1.59 to 1.15 V, SnO<sub>2</sub>: 1.79 to 1.15 V) because of increased recombination via more

available electron states (Q). Jsc is almost unchanged (24.1–24.9 mA/cm<sup>2</sup>), and FF achieves a plateau at moderate DOS (TiO<sub>2</sub>: 79.7%, WO<sub>3</sub>: 77.8%, SnO<sub>2</sub>: 78.7%) and slightly falls for very high DOS. Therefore, PCE lowers (TiO<sub>2</sub>: 28.0 to 22.5%, WO<sub>3</sub>: 26.2 to 20.9%, SnO<sub>2</sub>: 26.7 to 21.4%), with TiO<sub>2</sub> always performing better than SnO<sub>2</sub> and WO<sub>3</sub>, whereby low CB-DOS restricts recombination and achieves optimal efficiency[38].

Figure 12 illustrates As VB-DOS is raised from  $1 \times 10^{16}$  to  $1 \times 10^{20} \text{ cm}^{-3}$ , Voc decreases (TiO<sub>2</sub>: 3.93 to 1.21 V, WO<sub>3</sub>: 2.42 to 1.21 V, SnO<sub>2</sub>: 3.15 to 1.21 V) because of increased recombination from higher hole states (Q), while Jsc is little changed (23.7–24.9 mA/cm<sup>2</sup>)[39]. FF increases dramatically at moderate DOS (TiO<sub>2</sub>: 29.9 to 79.5%, WO<sub>3</sub> :

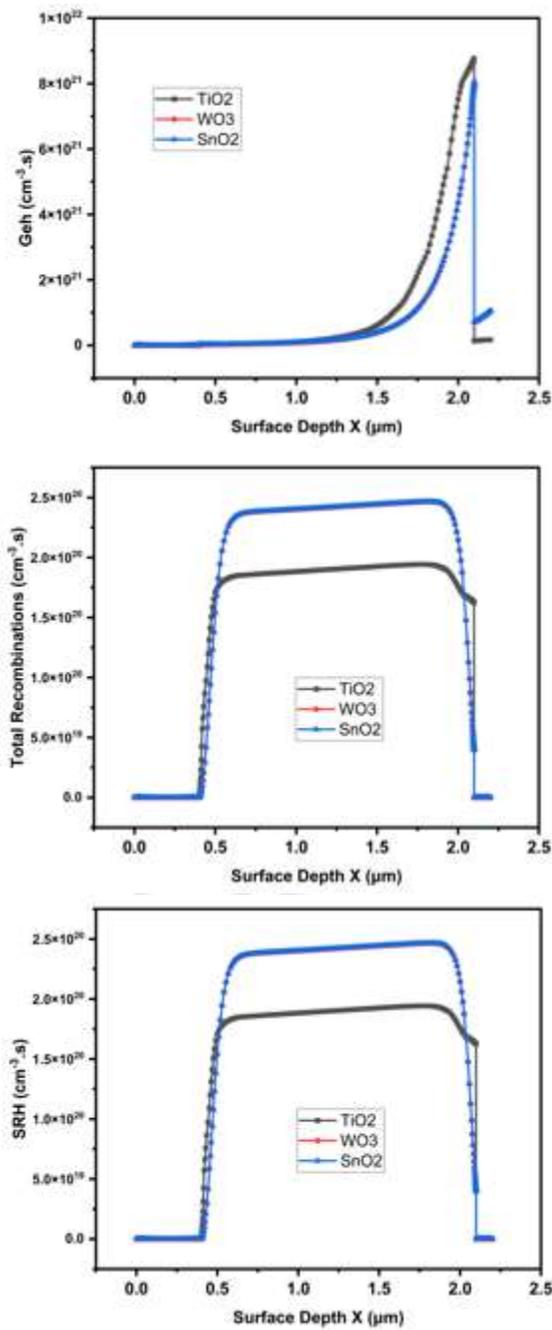


Figure 10: (a)-Generation , (b)-Total Recombination and (c)-SRH vs Surface depth of optimized solar cells with different ETMs.

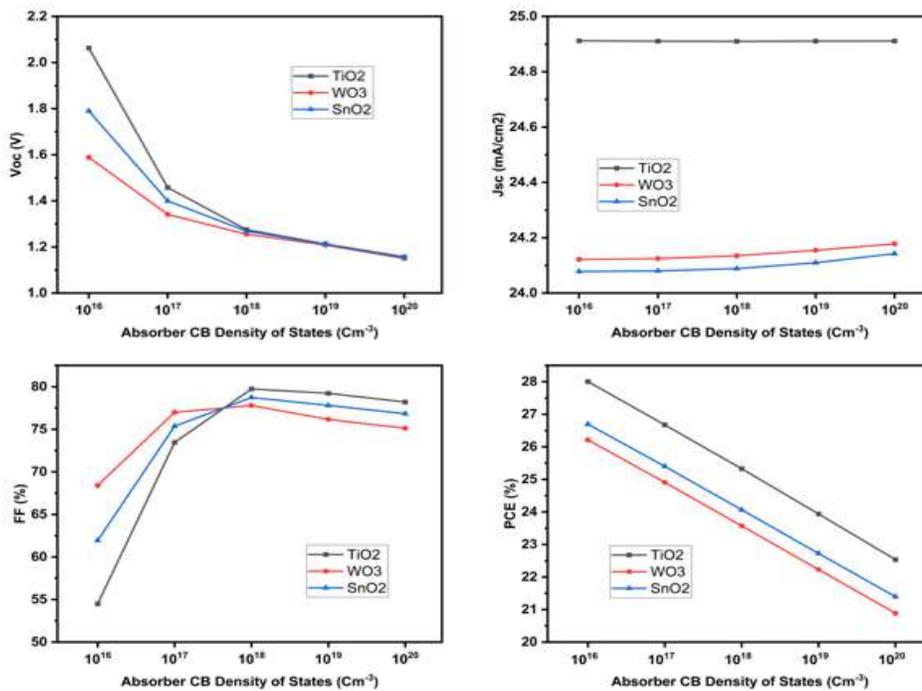


Figure 11: Variation of Electrical parameters vs Conduction band density of states (CB) of absorber layer with different ETMs

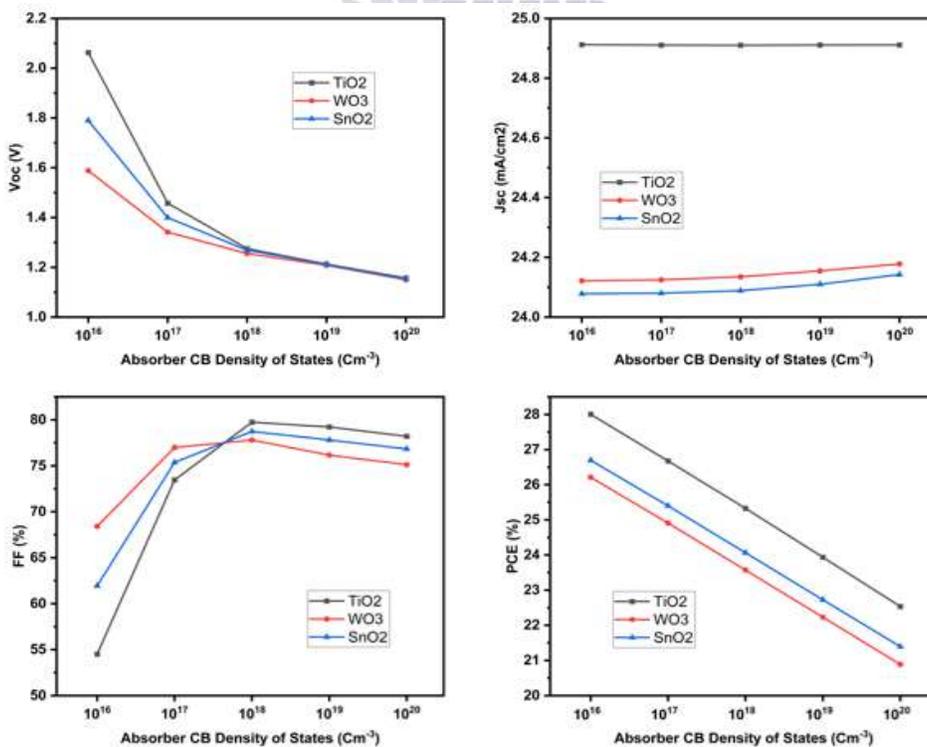


Figure 12: Variation of Electrical parameters vs Valance band density of states (VB) of absorber layer with different ETMs

47.3 to 77.8%, SnO<sub>2</sub>: 37.0 to 78.7%) before decreasing slightly at very high DOS. As a result, PCE drops with VB-DOS (TiO<sub>2</sub>: 29.2 to 23.9%, WO<sub>3</sub>: 27.2 to 22.2%, SnO<sub>2</sub>: 27.6 to 22.7%), with TiO<sub>2</sub> devices always outperforming SnO<sub>2</sub> and

WO<sub>3</sub>. This proves that low VB-DOS restricts hole-mediated recombination, maximizing efficiency, and that TiO<sub>2</sub> is the best ETM option for K<sub>2</sub>GeBr<sub>6</sub> solar cells.

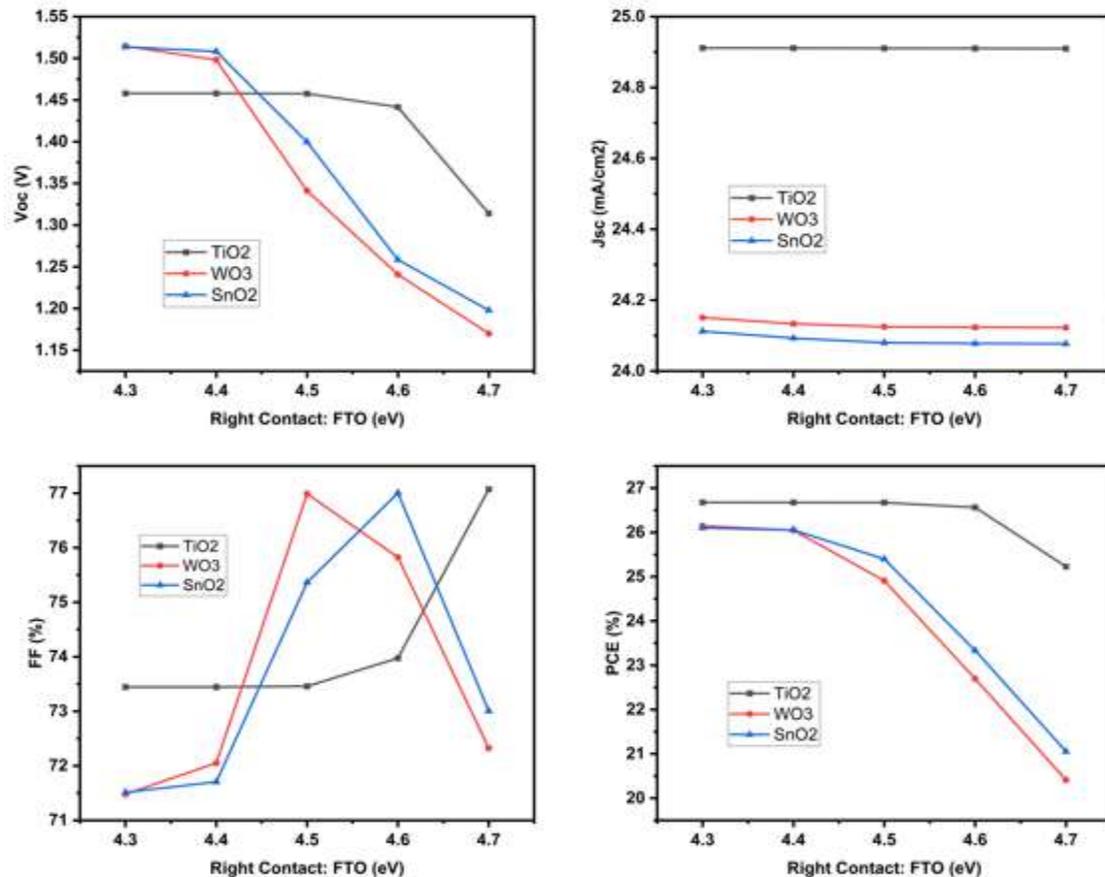


Figure 13: Variation of Electrical parameters vs Right contact (FTO) layer of optimized solar cells with different ETMs

The impact of the FTO work function (figure 13) on K<sub>2</sub>GeBr<sub>6</sub> devices with TiO<sub>2</sub>, WO<sub>3</sub>, and SnO<sub>2</sub> ETMs is presented in Figure 15. By increasing FTO work function from 4.3 to 4.7 eV, Voc goes down in all devices (TiO<sub>2</sub>: 1.458 to 1.314 V, WO<sub>3</sub>: 1.515 to 1.170 V, SnO<sub>2</sub>: 1.514 to 1.198 V) due to enhanced electron extraction energy barrier at the interface. Jsc is nearly constant (24.08–24.91 mA/cm<sup>2</sup>), whereas FF goes down slightly (TiO<sub>2</sub>: 73.45 to 77.07%, WO<sub>3</sub>: 71.48 to 72.33%, SnO<sub>2</sub>: 71.51 to 73.01%) based on charge carrier transport balance. Therefore, PCE comes down

(TiO<sub>2</sub>: 26.68 to 25.23%, WO<sub>3</sub>: 26.15 to 20.41%, SnO<sub>2</sub>: 26.10 to 21.06%), and TiO<sub>2</sub>-based cells outperform WO<sub>3</sub> and SnO<sub>2</sub> in all cases. This verifies that the best FTO work function is essential in order to minimize electron extraction barriers, suppress recombination at the contact, and achieve maximum efficiency[41]. Figure 14 illustrates the influence of Au metal work function on K<sub>2</sub>GeBr<sub>6</sub> devices with TiO<sub>2</sub>, WO<sub>3</sub>, and SnO<sub>2</sub> ETMs. With varying Au work function from 5.1 to 5.5 eV, Voc increases (TiO<sub>2</sub>: 1.458 to 1.920 V, WO<sub>3</sub>: 1.341 to 1.461 V, SnO<sub>2</sub>:

1.400 to 1.613 V) owing to better matching with the absorber, better hole extraction, and reduced interfacial recombination.  $J_{sc}$  is theoretically unaffected (24.08–24.91 mA/cm<sup>2</sup>), whereas FF decreases slightly in the case of TiO<sub>2</sub> (73.46 to 58.44%) but is excellent in the case of WO<sub>3</sub> (76.99 to 73.50%) and SnO<sub>2</sub> (75.36 to 67.96%) according to charge transport balance. PCE then

improves (TiO<sub>2</sub>: 26.67 to 27.95%, WO<sub>3</sub>: 24.91 to 25.91%, SnO<sub>2</sub>: 25.40 to 26.39%), with the maximum efficiency being delivered by TiO<sub>2</sub> devices followed by SnO<sub>2</sub> and WO<sub>3</sub>. This clearly shows that optimal Au work function improves hole extraction, reduces recombination at the back contact, and also maximizes device performance in general[42].

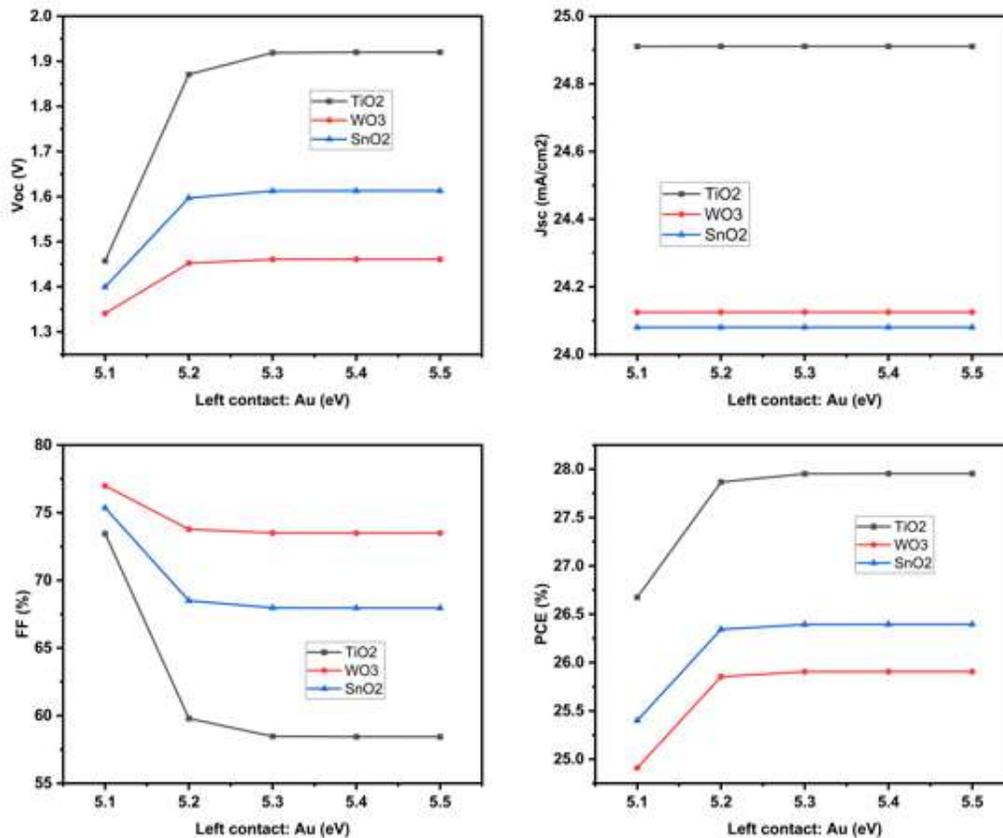


Figure 14: Variation of Electrical parameters vs Left contact (Au) layer of optimized solar cells with different ETMs

Figure 15 illustrates the influence of temperature (260–360 K) on K<sub>2</sub>GeBr<sub>6</sub> devices with TiO<sub>2</sub>, WO<sub>3</sub>, and SnO<sub>2</sub> ETMs. Voc decreases with increasing temperature (TiO<sub>2</sub>: 1.475 to 1.296 V, WO<sub>3</sub>: 1.249 to 1.294 V, SnO<sub>2</sub>: 1.294 to 1.295 V) because of increased carrier recombination and lower built-in potential at elevated temperatures.  $J_{sc}$  is comparatively stable (24.07–24.91 mA/cm<sup>2</sup>), while FF increases slightly (TiO<sub>2</sub>: 73.60 to 78.31%, WO<sub>3</sub>: 77.24 to 78.56%, SnO<sub>2</sub>: 78.56 to

78.55%) with decreased series resistance. PCE decreases accordingly (TiO<sub>2</sub>: 27.04 to 25.28%, WO<sub>3</sub>: 23.29 to 24.52%, SnO<sub>2</sub>: 24.52 to 24.49%), with TiO<sub>2</sub> devices being the most efficient at low temperatures, followed by SnO<sub>2</sub> and WO<sub>3</sub>. This indicates that higher temperature increases recombination, decreasing Voc and overall efficiency, but FF increases slightly due to lower resistance effects [43].

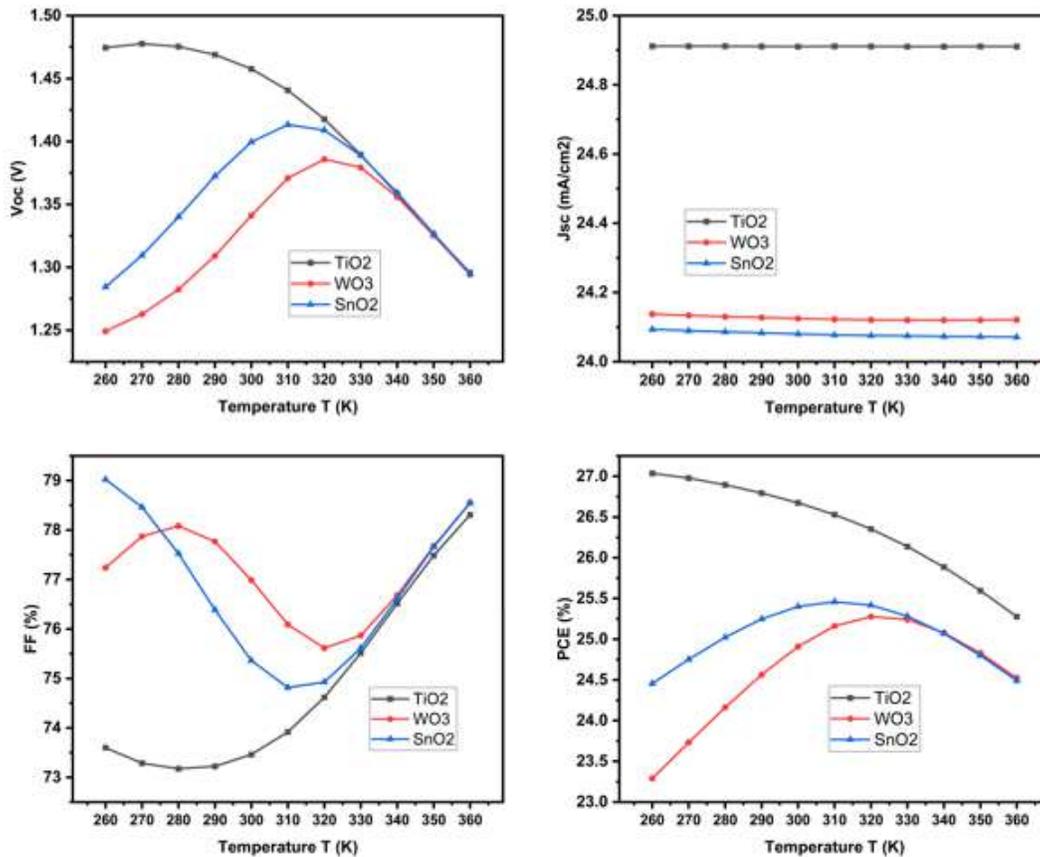


Figure 15: Variation of Electrical parameters vs Temperature of devices with different ETMs

Table 2: Solar cell device architecture vs optimized electrical parameters

Device Structure	Voc (V)	Jsc (mA/cm <sup>2</sup> )	FF (%)	PCE (%)
FTO/WoO <sub>3</sub> /K <sub>2</sub> GeBr <sub>6</sub> /TiO <sub>2</sub> /Au	1.4576	24.9106	73.46	26.67
FTO/WoO <sub>3</sub> /K <sub>2</sub> GeBr <sub>6</sub> /WO <sub>3</sub> /Au	1.3411	24.1247	76.99	24.91
FTO/WoO <sub>3</sub> /K <sub>2</sub> GeBr <sub>6</sub> /SnO <sub>2</sub> /Au	1.3996	24.0801	75.36	25.40

**Conclusion:**

Solar cell device level simulations of lead free K<sub>2</sub>GeBr<sub>6</sub> absorber with TiO<sub>2</sub>, WO<sub>3</sub>, and SnO<sub>2</sub> electron transport materials shows that TiO<sub>2</sub> is the best ETM with the highest carrier generation rate (10<sup>21</sup> cm<sup>-3</sup>s<sup>-1</sup>), lowest recombination loss (1.9×10<sup>20</sup> cm<sup>-3</sup>s<sup>-1</sup> as compared to 2.5×10<sup>20</sup> cm<sup>-3</sup>s<sup>-1</sup> for WO<sub>3</sub>/SnO<sub>2</sub>), and consistently better photovoltaic performance, with device electric

characteristics (PCE % ≈ 23–27%, Voc ≈ 1.3–1.46 V, Jsc ≈ 24 mA/cm<sup>2</sup>, FF ≈ 73–77%). The K<sub>2</sub>GeBr<sub>6</sub> is good absorbance performance, at an optimized thickness of 1 μm achieving a strong balance between the absorption of photons and recombination. The defect densities greater than 10<sup>16</sup> cm<sup>-3</sup> reduce efficiency, which underline the significance of interface and bulk defect

passivation. Overall, this study shows that  $\text{TiO}_2/\text{K}_2\text{GeBr}_6$  architecture as a high-efficiency and stable device combination, offering a sustainable alternative to lead-based perovskites and a strong foundation for the development of next-generation, eco-friendly solar cells.

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